Exponential Decay of Eigenfunctions of Partial Differential Operators

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### Introduction

In this talk I will consider the decay as  $|x| \to \infty$  of solutions  $\psi$  to the eigenvalue problem

$$(H - \lambda)\psi = 0.$$

Here H is in a class of self-adjoint partial differential operators of the form

$$H = Q(-i\nabla) + V(x)$$

where Q is a real elliptic polynomial and V is bounded. But first a review of some results for the Laplacian. Consider solutions  $\psi$  to

$$(-\Delta + V(x) - \lambda)\psi = 0, \psi \in L^2(\mathbb{R}^d),$$

 $V(x) \rightarrow 0$  as  $|x| \rightarrow \infty, \lambda$  and V real. Well known results: define  $\sigma_c$ 

$$\sigma_c = \sup\{\sigma : e^{\sigma r} \psi \in L^2(\mathbb{R}^d)\}$$

Theorem (O'Connor, Combes-Thomas, Bardos-Merigot, FH2HO)
i) If λ < 0 and V(x) = o(1) or λ > 0 and V(x) = o(|x|<sup>-1</sup>), then σ<sub>c</sub> > 0.
ii) If V(x) = O(|x|<sup>-1/2</sup>) then either σ<sub>c</sub> < ∞ or ψ = 0.</li>
iii) If V(x) = o(|x|<sup>-1/2</sup>) and σ<sub>c</sub> ∈ (0,∞) then λ < 0 and σ<sub>c</sub> = √-λ.
OPEN PROBLEM: If V is bounded or even V(x) = O(|x|<sup>-δ</sup>), δ > 0, is

it true that  $\sigma_c < \infty$  unless  $\psi = 0$ ?

#### Laplacian - the N-body problem

Consider now the N-body problem,  $H = -\Delta + V$ ,  $(H - \lambda)\psi = 0$  where for simplicity V is a sum of fast decaying 2-body potentials plus a fast decaying N-particle interaction.

Let

$$\Sigma_c = \{\sqrt{\tau - \lambda} : \tau \in T(H), \tau \ge \lambda\}$$

T(H) is the set of thresholds of H, a bounded closed countable set independent of the N-particle interaction.

Theorem (Froese, H)

Suppose  $H, \lambda, \psi (\neq 0)$  are as above. Then  $\sigma_c \in \Sigma_c$ .

The particular value of  $\sigma_c \in \Sigma_c$  will depend on the N-particle interaction as well as  $\lambda$ .

The point of this slide is to emphasize that there are situations (unlike the one-body problem with the Laplacian) where there may be several possibilities for the decay rate for a given  $\lambda$ .

## Results for $Q(-i\nabla)$

We consider  $L^2(\mathbb{R}^d)$  solutions  $\psi$  to

$$(Q(-i\nabla)+V(x)-\lambda)\psi=0$$

with Q a real elliptic polynomial and  $\lim_{|x|\to\infty} V(x) = 0$ . Consider the two conditions on a point  $(\sigma, \xi, \omega) \in \mathbb{R}_+ \times \mathbb{R}^d \times S^{d-1}$ 

$$Q(\xi + i\sigma\omega) = \lambda \tag{1}$$

$$P_{\perp}(\omega)\nabla_{\xi}Q(\xi+i\sigma\omega)=0$$
(2)

where  $P_{\perp}(\omega)$  is the projection onto the subspace of  $\mathbb{R}^d$  perpendicular to  $\omega$ .

#### Let

 $\Sigma_{c} = \{ \sigma > 0 : \exists (\xi, \omega) \in \mathbb{R}^{d} \times S^{d-1} \text{ so that } (1) \text{ and } (2) \text{ are satisfied} \}$ 

We will give conditions under which the decay rate  $\sigma_c$  is in  $\Sigma_c$ . Let  $V = V_1 + V_2$ , both real and bounded.

#### Theorem (There is exponential decay)

Under either of the following two conditions we can conclude that  $\sigma_c > 0$ :

**1**  $\lambda \notin RanQ$  and V(x) = o(1) at infinity

②  $\lambda \in RanQ$  but  $\lambda$  is not a critical value of Q and in addition  $\partial^{\alpha}V_1(x) = o(|x|^{-|\alpha|}) \forall \alpha, V_2(x) = o(|x|^{-1}).$  The next theorem eliminates the possibility of super-exponential decay (under strong conditions on the potential - too strong). Let q be the degree of the polynomial Q.

Theorem (There is no super-exponential decay)

Suppose  $V_2(x) = O(|x|^{-(q/2+\delta)})$ ,  $\partial^{\alpha} V_1(x) = O(|x|^{-(\delta+|\alpha|+q)/2})$ , where  $1 \le |\alpha| \le q$  and  $\delta > 0$ . Then  $\sigma_c < \infty$  or  $\psi = 0$ .

Given the last two theorems we know conditions on the potential so that  $\sigma_c \in (0, \infty)$ . In this case we can determine the decay rate "algebraically".

Theorem (Decay rate determined)

Assume  $\partial^{\alpha} V_1(x) = o(|x|^{-|\alpha|}) \forall \alpha, V_2(x) = o(|x|^{-1/2})$ . Suppose  $\sigma_c \in (0, \infty)$ . Then  $\sigma_c \in \Sigma_c$ .

# Example: $Q(\xi) = G(\xi^2)$

We take degree of Q = q. Then degree G = q/2. We assume  $\sigma \in (0,\infty)$ . With  $z = (\xi + i\sigma\omega)^2$  the critical equations (1) and (2) which determine the possible decay rates  $\sigma$  reduce to

 $G(z) = \lambda$  $G'(z)P_{\perp}(\omega)\xi = 0$ 

Except for at most (q - 2)/2 exceptional  $\lambda$ 's (arising from non-simple roots of  $G - \lambda$ ), these equations reduce to  $G((|\xi| + i\sigma)^2) = \lambda$ .

For an arbitrary real elliptic polynomial  $Q(\xi) = G(\xi^2)$ , every solution with positive  $\sigma$  of these equations occurs as a decay rate  $\sigma_c$  of an eigenfunction of  $G(-\Delta) + V$  for a real V in  $C_0^{\infty}(\mathbb{R}^d)$ .

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We only give some of the main ingredients of the proofs.

Condition (1) above,  $Q(\xi + i\sigma\omega) = \lambda$ , involves an energy estimate for the state  $e^{\sigma r}\psi$  which we do not discuss here.

To understand Condition (2),  $P_{\perp}(\omega)\nabla_{\xi}Q(\xi + i\sigma\omega) = 0$ , we use commutator methods with a special *conjugate operator* dependent on a parameter  $\sigma$  constructed as follows. First decompose  $Q(\xi + i\eta)$  into its real and imaginary parts:

$$Q(\xi + i\eta) = X(\xi, \eta) + iY(\xi, \eta)$$

Let

$$a(x,\xi)=rY(\xi,\sigma\omega(x))$$

Then the conjugate operator A is defined as the operator with Weyl symbol a:

$$A = Op^w(a)$$

Note that if  $Q(p) = p^2$ , then  $A = \sigma(x \cdot p + p \cdot x)$ ,  $p = -i\nabla$ . To get an idea where this comes from note that  $Q(p+i\sigma\omega) = e^{\sigma r}Q(p)e^{-\sigma r}$  and

$$i[Q(p), e^{\sigma r} A e^{\sigma r}] = e^{\sigma r} (i[\tilde{X}, A] + 2 \operatorname{Re} \tilde{Y} A) e^{\sigma r},$$

where  $\tilde{X} = \text{Re}(e^{\sigma r}Q(p)e^{-\sigma r})$  and  $\tilde{Y} = \text{Im}(e^{\sigma r}Q(p)e^{-\sigma r})$ . To leading order the symbol of the operator between exponentials to the right has symbol

$$r\{X,Y\} + 2rY^2 + \{X,r\}Y.$$
 (3)

We can calculate the Poisson bracket  $\{X, Y\}$  using the Cauchy - Riemann equations:

$$\{X,Y\} = |P_{\perp}(\omega) \nabla_{\xi} Q(\xi + i\sigma\omega)|^2$$

The last term in (3) can be bounded by the middle term and something of lower order.

#### Absence of super-exponential decay, $\sigma_c < \infty$

Two key points:  $\diamond$  No pseudo-differential operators. Exact computations.  $\diamond r = |x| \rightarrow r = \langle x \rangle - \langle x \rangle^{1-\epsilon} + 1$  (Rodnianski - Tao) Let  $\mathbf{a} = \mathbf{p} - i\sigma\omega, \omega = \nabla r$ . The eigenvalue equation for  $\psi_{\sigma} = e^{\sigma r}\psi$  is  $(Q(a^*) + V - \lambda)\psi_{\sigma} = 0$ 

Thus

 $egin{aligned} &<\psi_{\sigma}, ([\mathcal{Q}(a),\mathcal{Q}(a^{*})]+|\mathcal{Q}(a)+V_{1}-\lambda|^{2})\psi_{\sigma}> = \ &-<\psi_{\sigma}, (2 ext{Re}[\mathcal{Q}(a),V_{1}]+|V_{2}|^{2})\psi_{\sigma}> \end{aligned}$ 

Another key point:  $P := [a, a^*] \ge c\sigma r^{-(1+\epsilon)}$ 

We extract positivity from  $[Q(a), Q(a^*)]$ . Write the commutator as a sum of Wick ordered operators. Let  $J_m = (j_1, ..., j_m), K_m = (k_1, ..., k_m)$ . Then

$$[Q(a), Q(a^*)] = F + E$$
  

$$F = \sum_{m=1}^{q} \sum_{J_m, K_m} \partial_{J_m} Q(a^*) P_{J_m, K_m} \partial_{K_m} Q(a) / m!$$

Here  $\partial_{J_m} = \partial_{j_1} \cdots \partial_{j_m}$ ,  $\partial_{K_m} = \partial_{k_1} \cdots \partial_{k_m}$ ,  $P_{J_m,K_m} = P_{j_1,k_1} \cdots P_{j_m,k_m}$  and *E* is negligible for large  $\sigma$ .

Note for example that the term with m = q is bounded below by

 $c\sigma^q r^{-q(1+\epsilon)}$ .

Theorem (Improved super-exponential decay result)

*Take*  $Q(-i\nabla) = (-\Delta)^{q/2}, q = 2, 4.$ 

### Suppose

$$V_2(x) = O(|x|^{-q/4-\delta})$$
 $\partial^lpha V_1(x) = O(|x|^{-(\delta+q/2+|lpha|)/2})$  for  $1 \le |lpha| \le q/2$ , and  $\delta > 0$ .

#### Then $\sigma_c < \infty$ .

This theorem replaces q by q/2 in the assumptions on the potential for  $Q(-i\nabla) = (-\Delta)^{q/2}, q = 2, 4.$ 

#### Directional decay rates

First consider an **arbitrary**  $L^2$  function  $\phi$  with  $\sigma_c \in (0, \infty)$ . We introduce three exponential decay rates depending on a direction  $\omega \in S^{d-1}$ .

$$\sigma_{c}(\omega) = \sup\{\sigma | e^{\sigma \omega \cdot x} \phi \in L^{2}\}$$
  

$$\sigma_{k}(\omega) = \sup\{\eta \cdot \omega | \eta = \sigma_{c}(\nu)\nu \text{ for some } \nu \in S^{d-1}\}$$
  

$$\sigma_{loc}(\omega) = \sup\{\sigma | e^{\sigma |x|} \phi \in L^{2}(C) \text{ for some open cone C containing } \omega\}$$

We introduce the set

$$\mathcal{E} = \{\eta \in \mathbb{R}^d | e^{\eta \cdot x} \phi \in L^2\}$$

Note that  $\sigma_k(\omega) = \sup\{\eta \cdot \omega | \eta \in \mathcal{E}\}.$ 

#### Theorem

### • $\mathcal{E}$ is convex.

- **3**  $1/\sigma_c(\omega)$  is Lipschitz. In fact  $|1/\sigma_c(\omega_1)-1/\sigma_c(\omega_2)| \le |\omega_1-\omega_2|/\sigma_c$ .
- For  $\omega \in S^{d-1}$  define  $\sigma_k(r\omega) = r\sigma_k(\omega)$  for  $r \ge 0$ . Then  $\sigma_k(x)$  is the largest convex homogeneous function  $\lambda(x)$  of degree one which satisfies

$$e^{t\lambda(x)}\phi\in L^2$$
 for all  $t\in [0,1)$ .

If ρ is a continuous map of ℝ<sup>d</sup> → [0,∞) which is homogeneous of degree one satisfying ρ(ω) ≤ σ<sub>loc</sub>(ω) then

 $e^{t
ho(x)}\phi\in L^2$  for all  $t\in[0,1)$ .

#### Theorem

Suppose  $\phi$  is an eigenfunction of  $Q(-i\nabla) + V(x)$  with eigenvalue  $\lambda$  and with  $\sigma_c \in (0, \infty)$ . Suppose  $\partial^{\alpha} V_1(x) = o(|x|^{-\alpha}), V_2(x) = o(|x|^{-1/2})$ . Consider  $\omega$  such that  $\sigma_c(\omega) < \infty$ . Then for some  $(\xi, \theta, \beta) \in \mathbb{R}^d \times S^{d-1} \times \mathbb{C}, \sigma_c(\omega) = \sigma$  satisfies

 $Q(\xi + i\sigma\omega) = \lambda$  $\nabla Q(\xi + i\sigma\omega) = \beta\theta$ 

#### Theorem

Suppose  $\overline{\mathcal{E}}$  is strictly convex with a  $C^1$  boundary. Then  $\sigma_{loc}(\omega) = \sigma_k(\omega)$ , all  $\omega \in S^{d-1}$ .