EIGENVALUES AND APPROXIMATION THROUGH SIMPLE FUNCTIONS

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We consider two eigenvalue problems, associated with a continuous function and with his approximation through a simple (step) function. We prove that their corresponding eigenvalues have very different properties.

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1. INTRODUCTION

Consider $\mu_W < \mu_O$ and an increasing piecewise - C^1 function such that

(1)
$$\mu: (0,L) \to \mathbf{R}, \quad \mu_W \le \mu(x) \le \mu_O.$$

We study the eigenvalue problem

(2)
$$-(\mu f_x)_x + k^2 \mu f = \frac{1}{\sigma} k^2 f \mu_x, \quad x \in (0, L), \quad k > 0;$$

(3)
$$(\mu f_x f)^+(0) = \mu_W k f^2(0), \quad (\mu f_x f)^-(L) = -\mu_O k f^2(L);$$

where $^+$ and $^-$ are lateral limit values, k is a parameter (wave number) and σ, f are the eigenvalues and eigenfunctions.

We consider the same problem (2)-(3) when μ is approximated through a simple (step) function μ_S , with jumps in some interior points $x_i \in (0, L)$; then $(\mu_S)_x$ is a Dirac distribution. The corresponding eigenvalues and eigenfunctions are denoted by f_S, σ_S . The functions μ , μ_S are "very close" if the number of interior points x_i is very large (see the Figure 1 below).

When μ is linear and continuous, we get an upper bound of σ which is not depending on k. Moreover, we prove that σ becomes arbitrary small (positive) with increasing L. On the contrary, we get f_S such that the corresponding σ_S become infinite with increasing k, independent of L.

The main results of this paper is following. Even if the step function μ_S is very close to μ , there exists a strong difference between the eigenvalues corresponding to μ_S and μ .

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Graphs of a function f and the approximating simple function ϕ_1



Graphs of a function f and the approximating simple function ϕ_2

$$\phi_1(x) = \begin{cases} 0, & 0 \le f(x) < 1, \\ 1, & f(x) \ge 1. \end{cases}$$

$$\phi_2(x) = \begin{cases} 0, & 0 \le f(x) < \frac{1}{2}, \\ \frac{1}{2}, & \frac{1}{2} \le f(x) < 1, \\ 1, & 1 \le f(x) < \frac{3}{2}, \\ \frac{3}{2}, & \frac{3}{2} \le f(x) < 2, \\ 2, & f(x) \ge 2. \end{cases}$$



The system (1)-(2) is related with the linear stability of the displacements in a Hele-Shaw cell (a technical device described in [1], [10]). The three-layer Hele-Shaw model is studied in [2], [8], [9]. The papers [3]-[6] are concerning the multi-layer Hele - Shaw model. In these models, the boundary conditions (3) contain also σ . Therefore, even if the system (1)-(2) is much more simpler, we have a "singular behavior" of the eigenvalues. Some contradictions concerning the multi-layer model are proved in [11].

The very strong dependence of the eigenvalues of an algebraic matrix, as functions of the coefficients, has been highlighted in the well-known book [12].

2. WHEN μ IS A CONTINUOUS FUNCTION

We consider the linear continuous function μ such that

(4)
$$\mu(x) = (\mu_O - \mu_W)x/L + \mu_W, \quad \forall x \in [0, L].$$

We multiply with f in (2), we integrate on [0, L] and get

$$\int_{0}^{L} \mu f_{x}^{2} + k^{2} \int_{0}^{L} \mu f^{2} = \frac{k^{2}}{\sigma} \int_{0}^{L} \mu_{x} f^{2}.$$
and (5) give us

The relations (2)-(3) and (5) give us

(6)
$$\sigma_L = \frac{k^2 \int_0^L \mu_x f^2}{\mu_W k f^2(0) + \mu_O k f^2(L) + \int_0^L (\mu f_x^2 + k^2 \mu f^2)}.$$

We neglect some positive terms in the denominator and for the linear μ we obtain

(7)
$$\sigma_L \le \frac{\int_0^L \mu_x f^2}{\int_0^L \mu f^2} \le \frac{\mu_O - \mu_W}{L\mu_W}$$

In the general case, when $\mu(x)$ is an arbitrary continuous function on [0, L](not related with μ_W and μ_O), from (6) we obtain the following estimate for the corresponding eigenvalue, denoted by σ_C :

(8)
$$\sigma_C \le \frac{Max_x(\mu_x)}{Min_x(\mu)}.$$

Both above estimates are not depending on k. Moreover, the relation (7) gives us

(9)
$$L \to \infty \Rightarrow \sigma_L \to 0.$$

3. WHEN μ IS APPROXIMATED BY A SIMPLE (STEP) FUNCTION μ_S

We divide (0, L) in N small intervals, separated by the interfaces x_i ,

(10)
$$x_i = iL/N, \quad i = 0, 1, ..., N; \quad x_0 = 0, \quad x_N = L.$$

For i = 1, ..., N we introduce the step function

(11)
$$\mu_S(x) = \mu_W + i[\mu], \quad x \in [x_{i-1}, x_i); \quad [\mu] = [\mu_O - \mu_W]/(N+1).$$

When N = 3 we have

- three "intermediate" values μ_i , such that $\mu_W < \mu_1 < \mu_2 < \mu_3 < \mu_O$;

- four equally spaced interfaces $x_0 = 0$, $x_1 = L/3$, $x_2 = 2L/3$, $x_3 = L$ such that

$$\mu_S(x) = \mu_W + \mu_i, \ x \in [x_{i-1}, x_i);$$

- four equal jumps $\mu_1 - \mu_W = \mu_2 - \mu_1 = \mu_3 - \mu_2 = \mu_O - \mu_3 = (\mu_O - \mu_W)/4$. From now on, we omit the subscript *s*. The step function (11) is very "close" to the function (4) when N is large enough. The derivative μ_x is a Dirac distribution, therefore (see *Remark 2* below)

(12)
$$\int_0^L \mu_x f^2 = \sum_{i=1}^{i=N-1} f^2(x_i)[\mu]_i,$$

(13)
$$-\int_0^L (\mu f_x f)_x = (\mu^+ f_x^+ f)(x_0) - (\mu^- f_x^- f)(x_N) + \sum_{i=1}^{i=N-1} [\mu f_x]_i f(x_i),$$

where $[F]_i = F^+(x_i) - F^-(x_i)$. As μ_i is constant on each small interval, we get

(14)
$$-f_{xx} + k^2 f = 0, \quad x \in (x_{i-1}, x_i).$$

We consider the following particular solution of (14):

(15)
$$f(x) = \begin{cases} e^{kx} , x \in [x_0, x_{N-1}]; \\ f(x_{N-1})e^{-k(x-x_{N-1})}, x \in [x_{N-1}, x_N]. \end{cases}$$

The above function f is continuous, but f_x is not and (recall $x_N = L$)

(16)
$$f(x_{N-1}) = e^{kL(N-1)/N}, \quad f(x_N) = e^{k(2x_{N-1}-L)} = e^{kL(N-2)/N}.$$

We multiply (2) with f, then from (12), (13), (15) we get

(17)
$$\sigma = \frac{\sum_{i=1}^{i=N-1} k^2 [\mu]_i f_i^2}{DEN + \sum_{i=1}^{i=N} I_i},$$

(18)
$$DEN = \mu_1 k f_0^2 + \mu_N k f_N^2 + \sum_{i=1}^{i=N-2} [\mu]_i k f_i^2 - (\mu_{N-1} + \mu_N) k f_{N-1}^2,$$

(19)
$$I_i = \int_{x_{i-1}}^{x_i} \mu_i (f_x^2 + k^2 f^2), \quad f_i = f(x_i).$$

Remark 1. For large N we have DEN > 0. For this, we need (see (18))

(20)
$$\mu_1 + \mu_N f_N^2 + \sum_{i=1}^{i=N-2} [\mu]_i f_i^2 > (\mu_{N-1} + \mu_N) f_{N-1}^2.$$

We recall (16) and use the well known relation

$$y + y^{2} + \dots + y^{N-2} = (y^{N-1} - 1)/(y - 1), \quad y = e^{2kL/N},$$

therefore (20) is equivalent with

$$\mu_1 + (\mu_O - \frac{\mu_O - \mu_W}{N+1})e^{2kL(N-2)/N} + \frac{\mu_O - \mu_W}{N+1} \times \frac{y^{N-1}}{y-1} >$$

$$(\mu_O - 2\frac{\mu_O - \mu_W}{N+1})e^{2kL(N-1)/N}.$$

The last inequality holds for N large enough, when $(N-1)/N \to 1$, $(N-2)/N \to 1$.

By direct calculation we get

Lemma 1.

(21)
$$f(x) = e^{+kx}, \quad J(a,c) = \int_a^c (f_x^2 + k^2 f^2) \Rightarrow J(a,c) \le k \{ f^2(a) + f^2(c) \}.$$

The main result of this section is set out in the following proposition.

PROPOSITION 1.

(22)
$$\sigma \to \infty \quad when \quad k \to \infty.$$

Proof. From the relations (17), (20), (21), we get

(23)
$$\sigma \ge k \frac{\sum_{i=1}^{i=N-1} [\mu]_i e^{2kx_i}}{a_0 + \sum_{i=1}^{i=N-1} a_i e^{2kx_i} + a_N e^{2kz}}, \quad a_i > 0.$$

We prove that the maximum value of the exponential is the same in the numerator and the denominator of the above ratio. Indeed, from the relations (15) - (16) it follows

$$f(x_N) = e^{kz}, \quad z = \frac{2(N-1)L}{N} - L = \frac{L(N-2)}{N}.$$

Therefore

$$2x_{N-1} = 2\frac{(N-1)L}{N} > 2\frac{L(N-2)}{N} = 2z$$

It was very important that jumps $[\mu]_i$ were positive. \Box

Remark 2. We give a proof for the formula (12). Consider $\mu : [a, c] \to \mathbf{R}$, $b \in (a, c), \ \mu(x) = A, x \in [a, b); \quad \mu(x) = B, x \in [b, c]$. We see that $\mu_x(x) = 0$ for almost every $x \in [a, c]$. Only in the point x = b, we have

$$\mu_x(b) = \lim_{\epsilon \to 0} \frac{\mu(b) - \mu(b - \epsilon)}{\epsilon} = \lim_{\epsilon \to 0} \frac{B - A}{\epsilon}.$$

The next property is verified by a sufficiently smooth function F: there exists a point $\chi \in (b - \epsilon, b)$ such that $\int_{b-\epsilon}^{b} F(x) dx = \epsilon F(\chi)$. Therefore, when $\epsilon \to 0$ we get

$$\int_{a}^{c} \mu_{x}(x)F(x)dx = \int_{b-\epsilon}^{b} \mu_{x}(x)F(x)dx \to [\mu]_{b}F(b), \quad [\mu]_{b} = \mu^{+}(b) - \mu^{-}(b) = B - A$$

REFERENCES

- [1] J. Bear, Dynamics of Fluids in Porous Media. Elsevier, New York, 1972.
- [2] C. Carsso and G. Paşa, An optimal viscosity profile in the secondary oil recovery. RAIRO M2AN - Mod. Math. et Analyse Numérique, **32** (1998), 2, 211–221.
- [3] P. Daripa, Studies on stability in three-layer Hele-Shaw flows. Phys. Fluids., 20 (2008), 112101.
- [4] P. Daripa, Hydrodynamic stability of multi-layer Hele-Shaw flows. J. Stat. Mech., Art. No. P12005 (2008).
- [5] P. Daripa and X. Ding, Universal stability properties for Multi-layer Hele-Shaw flows and Applications to Instability Control. SIAM J. Appl. Math., 72 (2012), 1667–1685.
- [6] P. Daripa and X. Ding, A Numerical Study of Instability Control for the Design of an Optimal Policy of Enhanced Oil Recovery by Tertiary Displacement Processes. Transport in Porous Media 93 (2012), 675–703.
- [7] P. Daripa, Some Useful Upper Bounds for the Selection of Optimal Profiles. Physica A: Statistical Mechanics and its Applications **391** (2012). 4065–4069.
- [8] S. B. Gorell and G. M. Homsy, A theory of the optimal policy of oil recovery by secondary displacement process. SIAM J. Appl. Math. 43 (1983), 79–98.
- [9] S. B. Gorell and G. M. Homsy, A theory for the most stable variable viscosity profile in graded mobility displacement process. AIChE Journal, **31** (1985), 1598–1503.
- [10] H. S. Hele-Shaw, Investigations of the nature of surface resistence of water and of streamline motion under certain experimental conditions. Inst. Naval Architects Transactions 40(1898), 21–46.
- G. Paşa, A strong contradiction in multi-layer Hele-Shaw flow. arXiv:1903.01455 [physics.flu-dyn], Submitted on 4 Mar 2019.
- [12] J. H. Wilkinson, The algebraic eigenvalue problem. Oxford University Press, New York, 1988.

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