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EXISTENCE FOR A PARABOLIC EQUATION WITH
NONLINEAR BOUNDARY VALUE CONDITIONS

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1. INTRODUCTION

In this paper we are concerned with the nonlinear boundary-value problem of the form

$$\frac{\partial}{\partial t} \beta(u(t, x)) - \Delta u(t, x) \ni f(t, x), \quad (t, x) \in Q_T, \quad (1.1)$$

$$\frac{\partial u(t, x)}{\partial n} + \sigma(u(t, x)) \ni g(x, u(t, x)), \quad (t, x) \in \Sigma_T, \quad (1.2)$$

$$\beta(u(0, x)) \ni v_0(x), \quad x \in \Omega. \quad (1.3)$$

Here Ω is a bounded and open subset of the Euclidean space R^m with the boundary S ; $\frac{\partial u}{\partial n}$ is the outward normal derivative of u ; β and σ are maximal monotone graphs in R^2 (possible multi-valued); f, g and v_0 are given functions on $Q_T =]0, T[\times \Omega, S \times R$ and Ω respectively; $\Sigma_T =]0, T[\times S$.

Problems of this type occur in the heat radiation (see e.g. H.B. Keller [1]), the absorption of gas in a liquid (see e.g. C.V. Pao [2]), the thermostat control problems (see e.g. G. Duvaut and J.L. Lions [3]) and in description of other physical problems. For instance, if we consider the following boundary control problem governed by the Stefan problem (see Ch. Saguez [4])

$$\frac{\partial G(\theta)}{\partial t} - \Delta \theta = 0, \quad (t, x) \in Q_T,$$

$$\frac{\partial \theta}{\partial n} + h(\theta - \theta_e) = u, \quad (t, x) \in \Sigma_T,$$

$$G(\theta(0, x)) \ni v_0(x), \quad x \in \Omega,$$

where G is a maximal monotone graph in R^2 , h is a positive constant, θ_e

and v_0 are given functions, the existence result which will be given below allows to implement a feedback law of the following form

$$u(t, x) = g(x, u(t, x))$$

where g satisfies conditions $(g)_1, (g)_2$ given below. For other results concerning problem (1.1) - (1.3) in the case $g \equiv 0$ we refer the reader to the works of J.L. Lions [5] and Ph. Benilan [6].

We shall suppose here that the boundary S of Ω is sectionally smooth, i.e. it consists of k disjoint parts S_i , $i = 1, \dots, k$, such that $S = \bigcup_{i=1}^k \bar{S}_i$ and $S_i \cap S_j = \emptyset$, $1 \leq i < j \leq k$, and there exists for each S_i a cartesian coordinate system $y = T_i x$ such that

$$S_i : y_n = h_i(\bar{y}), \quad \bar{y} = (y_1, \dots, y_{m-1}) \in D_i, \quad (1.4)$$

$$\Omega_\lambda^i = \left\{ T_i^{-1} y; \bar{y} \in D_i, 0 < y_m - h_i(\bar{y}) < \lambda^{\frac{1}{2}} \right\} \subset \Omega, \quad (1.5)$$

where D_i is a bounded and open subset of R^{m-1} ; h_i is a once continuously differentiable function in D_i , and its gradient is bounded on D_i ; λ is positive and sufficiently small.

As regards the maximal monotone graphs β and σ and functions f , g and v_0 , we shall suppose that

$$(\sigma) \quad 0 \in \sigma(0), \quad D(\sigma) = R^1;$$

$$(\beta)_1 \quad 0 \in \beta(0), \quad R(\beta) = R^1 \quad \text{and} \quad \text{int } D(\beta) \ni 0;$$

$$(\beta)_2 \quad \text{for each } N > 0 \text{ there exists } \alpha_N > 0 \text{ such that}$$

$$(\beta u - \beta v)(u - v) \geq \alpha_N (u - v)^2$$

for $u, v \in D(\beta)$ with $|u|, |v| \leq N$;

$(g)_1$ $g(x, u)$ is measurable as a function of $x \in S$ for each $u \in R^1$, and continuous as a function of $u \in R^1$ for a.e. $x \in S$. In addition, it is assumed that for each $N > 0$ there exist $M_N > 0$, $L_N > 0$ such that

$$|g(x, u)| \leq M_N,$$

$$|g(x, u) - g(x, v)| \leq L_N |u - v|$$

for a.e. $x \in S$ and $u, v \in R^1$ with $|u|, |v| \leq N$;

$(g)_2$ there exist $k_1 < 0$, $k_2 > 0$ such that $[k_1, k_2] \subset \text{int } D(\beta)$ and $g(x, u)u \leq 0$ for a.e. $x \in S$ and $u \notin [k_1, k_2]$;

$$(f) \quad f \in L^2(0, T; L^\infty(\Omega));$$

$$(v_0) \quad v_0 \in H^1(\Omega) \cap L^\infty(\Omega),$$

where $H^1(\Omega)$ denotes the usual Sobolev space $W^{1,2}(\Omega)$

We shall denote by $H = L^2(\Omega)$. Let $W' = (H^1(\Omega))'$ be the dual space of $W = H^1(\Omega)$, i.e. W' is the completion of H under the norm

$$\|w^*\|_{W'} = \sup_{\|u\|_W=1} (w^*, u). \quad (1.6)$$

Here $(\cdot, \cdot)$ is the usual inner product in H and will be also used to denote the pairing between W and W' . The norms in H, W and W' are denoted by $\|\cdot\|_H, \|\cdot\|_W$ and $\|\cdot\|_{W'}$, respectively. As usual, $H^1(Q_T)$ denotes the Sobolev space $W^{1,2}(Q_T)$. Let X be a Banach space, we shall denote by $C([0, T]; X)$ the space of X -value continuous functions on $[0, T]$. $L^p(0, T; X)$, $p \geq 1$, denotes the space of functions $t \rightarrow h(t)$ measurable from $[0, T]$ to X (for the measure dt) such that

$$\left(\int_0^T \|h(t)\|_X^p dt \right)^{\frac{1}{p}} = \|h\|_{L^p(0, T; X)} < +\infty, \quad (p \neq +\infty)$$

$$\operatorname{ess\,sup}_{t \in]0, T[} \|h(t)\|_X = \|h\|_{L^\infty(0, T; X)} < +\infty, \quad (p = +\infty)$$

Let $A: W \rightarrow W'$ be the operator defined by

$$Au = \left\{ w^* \in W'; \quad \exists z \in L^2(S), z(x) \in \sigma(u(x)) \text{ a.e. } x \in S \text{ such that} \right. \\ \left. (w^*, \psi) = \int_\Omega \operatorname{grad} u \operatorname{grad} \psi \, dx + \int_S z \psi \, ds, \forall \psi \in W \right\} \quad (1.7)$$

for $u \in D(A) = \{u \in W; Au \neq \emptyset\}$.

For $u \in W$ such that $g(\cdot, u(\cdot)) \in L^2(S)$, we set $Gu \in W'$ defined by

$$(Gu, \psi) = \int_S g(x, u(x)) \psi(x) \, ds, \quad \forall \psi \in W. \quad (1.8)$$

Let $B: H \rightarrow H$ be the operator defined by

$$Bu = \left\{ v \in H; \quad v(x) \in \beta(u(x)) \text{ a.e. } x \in \Omega \right\} \quad (1.9)$$

for $u \in D(B) = \{u \in H; Bu \neq \emptyset\}$. As it is well known, B is a maximal monotone operator on H .

Therefore, problem (1.1) - (1.3) can be written under the following form

$$(Bu(t))' + Au(t) - Gu(t) \ni f(t), \quad t \in]0, T[\quad (1.10)$$

$$Bu(0) \ni v_0, \quad (1.11)$$

where $' = \frac{d}{dt}$.

The main result of this paper is the following theorem.

THEOREM 1.1 Let conditions $(\sigma), (\beta)_{1-2}, (g)_{1-2}, (f)$ and (v_0) be satisfied. Then problem (1.1)-(1.3) has a solution $u \in C([0, T]; H) \cap L^\infty(0, T; W) \cap L^\infty(Q_T)$ in the sense that there exists $v \in L^2(0, T; H)$ with $v(t) \in Bu(t)$ a.e. $t \in]0, T[$ such that

$$\frac{d}{dt}(v(t), \psi) + (Au(t) - gu(t), \psi) = (f(t), \psi), \quad \text{a.e. } t \in]0, T[, \quad (1.12)$$

$$v(0) = v_0 \quad (1.13)$$

for all $\psi \in W$.

The contents of this paper is outlined below. In Section 2 we shall construct an approximating problem associated with (1.1)-(1.3). We shall prove the existence and we shall establish a priori estimates for its solutions. The proof of the main theorem is delivered in Section 3.

2. APPROXIMATING PROBLEM

Let A_H be the operator $H \rightarrow H$ defined by $A_H u = Au \cap H$ for $u \in D(A_H) = \{u \in W; Au \cap H \neq \emptyset\}$. It is easy to show that $w^* \in A_H u$, if and only if u is a solution in $H^1(\Omega)$ to the problem

$$-\Delta u = w^*, \quad x \in \Omega, \quad (2.1)$$

$$\frac{\partial u}{\partial n} + \sigma(u) \geq 0, \quad x \in S. \quad (2.2)$$

We shall denote by $j: R^1 \rightarrow R^1$ a convex continuous function such that $\sigma = \partial j$ and $j \geq 0$ on R^1 , where ∂j denotes the subdifferential of j . It is well known that such a function always exists (see, e.g., [7, pp. 59-60]).

Therefore, $A_H = -\Delta$, where $D(A_H) = \{u \in H^2(\Omega); -\frac{\partial u}{\partial n} \in \sigma(u) \text{ a.e. } x \in S\}$, and $A_H = \partial \varphi$, where φ is the proper convex lower-semicontinuous function from H to $]-\infty, +\infty]$ defined by

$$\varphi(u) = \begin{cases} \frac{1}{2} \int_{\Omega} |\text{grad } u|^2 dx + \int_S j(u) dS, & \text{if } u \in H^1(\Omega) \text{ and } j(u) \in L^1(S), \\ +\infty, & \text{otherwise.} \end{cases} \quad (2.3)$$

(see [7, pp. 63-67]).

Let $A_\lambda = A_H(1 + \lambda A_H)^{-1}$ be the Yosida approximation of A_H .

Next we shall approximate the operator G defined by (1.8) by a family of Lipschitz continuous operators on H . To make whole what follows meaningful, let us briefly describe the physical situation from which this approximation originates.

It is well known that if problem (1.1)-(1.3) describes a diffusion process of heat in a domain Ω , $u(t, x)$ denoting the temperature of a point $x \in \Omega$ at time t , then the function g is a part of the surface density of the heat injection through the boundary S of Ω into its interior. In other words, the heat injection, which is given by surface density g in unity time through an element ΔS of surface S , is equal to integral $\int_{\Delta S} g ds$. The function f is the volume density of the heat injection directly in interior of Ω . In other words, the heat injection, which is given by volume density f in unity time through an element $\Delta \Omega$ of volume Ω , is equal to integral $\int_{\Delta \Omega} f dx$. Let $\hat{g}$ be a function defined on Ω and satisfying the following conditions

$$\text{Supp } \hat{g} \subset \sum_{i=1}^k \Omega_\lambda^i, \quad (2.4)$$

$$\int_{\Delta S} g ds = \sum_{i=1}^k \int_{\Omega_\lambda^i} g dx, \quad (2.5)$$

where $\lambda > 0$ is sufficiently small and

$$\Delta \Omega_\lambda^i = \left\{ x \in \Omega_\lambda^i ; \bar{x}^i \in \Delta S \cap S_i \right\}, \quad (2.6)$$

where $\bar{x}^i$ is given by

$$\bar{x}^i = T_i^{-1}(\overline{T_i x}, h_i(\overline{T_i x})) \quad (2.7)$$

i.e., the projection of x on S_i in the direction of y_m -axis.

Then it is physically obvious that the problem obtained by substituting f and g by 0 and $f + \hat{g}$ respectively, is a natural approximation to original problem.

To this aim, we set

$$g_{\lambda, i}(x, u) = \begin{cases} \lambda^{-\frac{1}{2}} g(\bar{x}^i, u) n_i(x), & x \in \Omega_\lambda^i, u \in R^1, \\ 0, & \text{otherwise} \end{cases} \quad (2.8)$$

$$g_{\lambda}(x, u) = \sum_{i=1}^k g_{\lambda, i}(x, u), \quad x \in \Omega, u \in \mathbb{R}^1, \quad (2.9)$$

where

$$n_i(x) = \left[1 + \sum_{j=1}^{m-1} \left(\frac{\partial h_j}{\partial y_j}(\bar{T}_i x) \right)^2 \right]^{\frac{1}{2}} \quad (2.10)$$

We define the operator $G_{\lambda} : H \rightarrow H$ by

$$(G_{\lambda} u)(x) = g_{\lambda}(x, u(x)), \quad \text{a.e. } x \in \Omega \quad (2.11)$$

for $u \in H$ such that $g_{\lambda}(\cdot, u(\cdot)) \in H$.

Thus, for each $\lambda > 0$, we consider the approximating problem

$$(Bu_{\lambda}(t))' + A_{\lambda} u_{\lambda}(t) - G_{\lambda} u_{\lambda}(t) \ni f(t), \quad t \in]0, T[, \quad (2.12)$$

$$Bu_{\lambda}(0) \ni v_0. \quad (2.13)$$

We have the following lemma.

LEMMA 2.1. Let assumptions $(\sigma), (\beta)_{1-2}, (g)_{1-2}, (f)$ and (v_0) be satisfied. Then, problem (2.12) - (2.13) for $q \geq 2$ has a unique solution

$u_{\lambda} \in C([0, T]; L^q(\Omega)) \cap L^{\infty}(Q_T)$ with $u_{\lambda}' \in L^{\infty}(0, T; L^q(\Omega))$ in the sense that there exists $v_{\lambda} \in C([0, T]; L^q(\Omega)) \cap L^{\infty}(Q_T)$ with $v_{\lambda}' \in L^{\infty}(0, T; L^q(\Omega))$ such that $v_{\lambda}(t) \in Bu_{\lambda}(t)$ a.e. $t \in]0, T[$ and

$$v_{\lambda}'(t) + A_{\lambda} u_{\lambda}(t) - G_{\lambda} u_{\lambda}(t) = f(t), \quad t \in]0, T[, \quad (2.14)$$

$$v_{\lambda}(0) = v_0. \quad (2.15)$$

PROOF. We define

$$g^N(x, u) = \begin{cases} g(x, u), & x \in S, |u| \leq N, \\ g(x, N), & x \in S, u > N, \\ g(x, -N), & x \in S, u < -N, \end{cases} \quad (2.16)$$

$$g_{\lambda, i}^N(x, u) = \begin{cases} \lambda^{-\frac{1}{2}} g^N(\bar{x}^i, u) n_i(x), & x \in \Omega_{\lambda}^i, u \in \mathbb{R}^1, \\ 0, & \text{otherwise,} \end{cases} \quad (2.17)$$

$$g_{\lambda}^N(x, u) = \sum_{i=1}^k g_{\lambda, i}^N(x, u), \quad x \in \Omega, u \in \mathbb{R}^1 \quad (2.18)$$

and the operator $G_{\lambda}^N : H \rightarrow H$ by

$$(G_{\lambda}^N u)(x) = g_{\lambda}^N(x, u(x)), \quad \text{a.e. } x \in \Omega \quad (2.19)$$

for each $u \in H$.

Let $\beta_N = \beta + \partial I_N$, where

$$I_N(u) = \begin{cases} 0, & |u| \leq N, \\ +\infty, & \text{otherwise.} \end{cases}$$

Then, β_N is a maximal monotone graph in R^2 , because $0 \in \text{int } D(\partial I_N) \cap D(\beta) \neq \emptyset$ (see [3, p.46]). We define $B_N: H \rightarrow H$ by

$$B_N u = \left\{ v \in H; v(x) \in \beta_N(u(x)) \text{ , a.e. } x \in \Omega \right\} \quad (2.20)$$

for $u \in D(B_N) = \left\{ u \in H; B_N u \neq \emptyset \right\}$.

First, we consider the problem

$$(B_N u_\lambda^N(t))' + A_\lambda u_\lambda^N(t) - G_\lambda^N u_\lambda^N(t) \ni f(t), \quad t \in]0, T[, \quad (2.21)$$

$$B_N u_\lambda^N(0) \ni v_0, \quad (2.22)$$

or equivalently

$$u_\lambda^N(t) = B_N^{-1} \left\{ v_0 + \int_0^t (f(\tau) - A_\lambda u_\lambda^N(\tau) + G_\lambda^N u_\lambda^N(\tau)) d\tau \right\}. \quad (2.23)$$

Since for $u \in L^q(\Omega)$, $q \geq 2$, the problem

$$y(x) - \lambda \Delta y(x) = u(x), \quad x \in \Omega, \quad (2.24)$$

$$\frac{\partial y(x)}{\partial n} + \sigma(y(x)) \ni 0, \quad x \in S \quad (2.25)$$

has a unique solution in $H^2(\Omega) \cap L^q(\Omega)$ (see H. Brézis [8, pp.58-59]), it follows that $(1 + \lambda A_H)^{-1}$ is nonexpansive on $L^q(\Omega)$, and therefore, A_λ is Lipschitz on $L^q(\Omega)$. Noticing that B_N^{-1} , G_λ^N are also Lipschitz on $L^q(\Omega)$, it follows by standard arguments that problem (2.23) has a unique solution $u_\lambda^N \in C([0, T]; L^q(\Omega))$ with $(u_\lambda^N)' \in L^\infty(0, T; L^q(\Omega))$.

To conclude the proof of Lemma 2.1, we shall show that for sufficiently large N ,

$$\|u^N(t)\|_{L^\infty(\Omega)} < N \text{ for all } t \in [0, T] \text{ and } \lambda > 0. \quad (2.26)$$

Let

$$v_\lambda^N(t) = v_0 + \int_0^t (f(\tau) - A_\lambda u_\lambda^N(\tau) + G_\lambda^N u_\lambda^N(\tau)) d\tau. \quad (2.27)$$

Then $v_\lambda^N \in B_N u^N$, $v_\lambda^N \in C([0, T]; L^q(\Omega))$ and

$$(v_\lambda^N(t))' + A_\lambda u_\lambda^N(t) - G_\lambda^N u_\lambda^N(t) = f(t), \text{ a.e. } t \in]0, T[, \quad (2.28)$$

$$v_\lambda^N(0) = v_0. \quad (2.29)$$

It is easy to show that

$$\int_0^t \int_\Omega |v_\lambda^N(\tau)|^{q-2} v_\lambda^N(\tau) (v_\lambda^N(\tau))' dx d\tau = \frac{1}{q} \|v_\lambda^N(t)\|_{L^q(\Omega)}^q - \frac{1}{q} \|v_0\|_{L^q(\Omega)}^q, \quad (2.30)$$

$$\int_0^t \int_\Omega f(\tau) |v_\lambda^N(\tau)|^{q-2} v_\lambda^N(\tau) dx d\tau \leq \int_0^t \|f(\tau)\|_{L^q(\Omega)} \|v_\lambda^N(\tau)\|_{L^q(\Omega)}^{q-1} d\tau. \quad (2.31)$$

Using PROPOSITION 1.1 in [7, p.183], we may infer that

$$\int_0^t \int_{\Omega} A_{\lambda} u_{\lambda}^N(\tau) |v_{\lambda}^N(t)|^{q-2} v_{\lambda}^N(\tau) dx d\tau \geq 0. \quad (2.32)$$

By virtue of $(\beta)_1, (g)_2$ and taking $N > \max \{-k_1, k_2\}$, we deduce

$$\int_0^t \int_{\Omega} g_{\lambda}^N(x, u_{\lambda}^N(\tau, x)) |v_{\lambda}^N(\tau, x)|^{q-2} v_{\lambda}^N(\tau, x) dx d\tau \leq \quad (2.33)$$

$$\leq \sum_{i=1}^k \lambda^{-\frac{1}{2}} t \operatorname{mes} \Omega_{\lambda}^i \cdot C_g \sup_{x \in \Omega_{\lambda}^i} \eta_i(x) \cdot C_{\beta}^{q-1},$$

where $C_g = \operatorname{ess\,sup}_{(x,u) \in S \times [k_1, k_2]} g(x, u)$, and $C_{\beta} = \sup_{u \in [k_1, k_2]} |\beta(u)|$.

Multiplying (2.28) by $|v_{\lambda}^N(t)|^{q-2} v_{\lambda}^N(t)$, integrating over Q_t , and using (2.30) - (2.33), we get

$$\|v_{\lambda}^N(t)\|_{L^q(\Omega)}^q \leq (\|v_0\|_{L^q(\Omega)}^q + C_q C_{\beta}^{q-1}) + q \int_0^t \|f(\tau)\|_{L^q(\Omega)} (1 + \|v_{\lambda}^N(\tau)\|_{L^q(\Omega)}^q) d\tau. \quad (2.34)$$

According to Gronwall's inequality, we obtain

$$\begin{aligned} \|v_{\lambda}^N(t)\|_{L^q(\Omega)}^q &\leq (\|v_0\|_{L^q(\Omega)}^q + C_q C_{\beta}^{q-1} + q \|f\|_{L^1(0,t;L^q(\Omega))}) \cdot \\ &\cdot (1 + q \|f\|_{L^1(0,t;L^q(\Omega))}) e^{q \|f\|_{L^1(0,t;L^q(\Omega))}}. \end{aligned} \quad (2.35)$$

Taking q -root and letting $q \rightarrow +\infty$, it follows that

$$\|v_{\lambda}^N(t)\|_{L^{\infty}(\Omega)} \leq b, \quad (2.36)$$

for all $t \in [0, T]$, where $b = (1 + C_{\beta} + \|v_0\|_{L^{\infty}(\Omega)}) (1 + e^{\|f\|_{L^1(0,T;L^{\infty}(\Omega))}})$

is independent of N , λ and t .

From $(\beta)_{1-2}$, we see that β^{-1} is locally Lipschitz on $\mathbb{R}^1$. Noticing that $|\beta^{-1}v| \geq |\beta_N^{-1}v|$ for all $v \in \mathbb{R}^1$, we see by (3.36) that

$$\|u_{\lambda}^N(t)\|_{L^{\infty}(\Omega)} \leq \max \{ \beta^{-1}(b), -\beta^{-1}(-b) \} \quad (2.37)$$

for all $t \in [0, T]$.

Taking $N > \max \{-k_1, k_2, \beta^{-1}(b), -\beta^{-1}(-b)\}$, we obtain (2.26) as claimed. For this N , letting $v_{\lambda} = v_{\lambda}^N$, we see that $u_{\lambda} = u_{\lambda}^N$ is a solution of problem (2.12)-(2.13). This completes the proof of Lemma 2.1.

In order to establish a priori estimates for the solution u_{λ} given by Lemma 2.1, we need of the following lemma.

LEMMA 2.2 If $u \in W \cap L^{\infty}(\Omega)$, then $u \in L^{\infty}(S)$ and $\|u\|_{L^{\infty}(S)} \leq \|u\|_{L^{\infty}(\Omega)}$.

PROOF. For $x \in \Omega_{\lambda}^1$, we denote (for simplicity)

$$u(y) = u(\bar{y}, y_m) = u(T_i^{-1} y) = u(x)$$

where $y = T_i x$. For $q \geq 2$ we have

$$\begin{aligned} |u(\bar{y}, h_i(\bar{y}))|^q &= |u(y)|^q + q \int_{y_m}^{h_i(\bar{y})} |u(\bar{y}, \xi)|^{q-2} u(\bar{y}, \xi) \frac{\partial u(\bar{y}, \xi)}{\partial \xi} d\xi \\ &\leq |u(y)|^q + q \int_{h_i(\bar{y})}^{h_i(\bar{y}) + \sqrt{\lambda}} |u(\bar{y}, \xi)|^{q-1} \left| \frac{\partial u(\bar{y}, \xi)}{\partial \xi} \right| d\xi, \end{aligned}$$

and therefore,

$$|u(\bar{y}, h_i(\bar{y}))|^q \leq \int_{h_i(\bar{y})}^{h_i(\bar{y}) + \sqrt{\lambda}} \left(\lambda^{-\frac{1}{2}} |u(y)|^q + q |u(y)|^{q-1} \left| \frac{\partial u(y)}{\partial y_m} \right| \right) dy_m. \quad (2.38)$$

Multiplying (2.38) by $\eta_i(T_i^{-1} y)$, integrating over D_i and using the boundness of η_i , it follows that

$$\|u\|_{L^q(S_i)}^q \leq c \left(\lambda^{-\frac{1}{2}} \|u\|_{L^q(\Omega_{\lambda}^1)}^q + q \|u\|_{L^{2(q-1)}(\Omega_{\lambda}^1)}^{q-1} \left\| \frac{\partial u}{\partial y_m} \right\|_{L^2(\Omega_{\lambda}^1)} \right).$$

Using the inequality

$$\|u\|_{L^q(\Omega_{\lambda}^1)}^q \leq \|u\|_{L^{2(q-1)}(\Omega_{\lambda}^1)}^q (\text{mes } \Omega_{\lambda}^1)^{1 - \frac{q}{2(q-1)}}$$

and taking q -root, we infer

$$\begin{aligned} \|u\|_{L^q(S_i)} &\leq c^{\frac{1}{q}} \left[\lambda^{-\frac{1}{2}} (\text{mes } \Omega_{\lambda}^1)^{1 - \frac{q}{2(q-1)}} \|u\|_{L^{2(q-1)}(\Omega_{\lambda}^1)} + \right. \\ &\quad \left. + q \left\| \frac{\partial u}{\partial y_m} \right\|_{L^2(\Omega_{\lambda}^1)} \right]^{\frac{1}{q}} \|u\|_{L^{2(q-1)}(\Omega_{\lambda}^1)}^{1 - \frac{1}{q}}. \quad (2.39) \end{aligned}$$

Letting $q \rightarrow +\infty$ in (2.39), we complete the proof of Lemma 2.2.

LEMMA 2.3 The solution u_{λ} , given by Lemma 2.1, satisfies the following estimates:

$$\lambda |A_H y_{\lambda}(t)|_H^2 + \|y_{\lambda}(t)\|_W^2 + \int_0^t |u'_{\lambda}(\tau)|_H^2 d\tau \leq c, \quad (2.40)$$

$$\int_0^t |y'_{\lambda}(\tau)|_H^2 d\tau \leq c, \quad (2.41)$$

$$\|y_{\lambda}(t)\|_{L^{\infty}(S)} \leq c, \quad (2.42)$$

$$\|A_H y_\lambda(t)\|_{W_1} \leq C, \quad (2.43)$$

where $y_\lambda(t) = (1 + \lambda A_H)^{-1} u_\lambda(t)$ and C is a constant independent of $\lambda > 0$ and $t \in [0, T]$.

PROOF. Multiplying (2.14) by $u'_\lambda(t)$, integrating over Q_t and using the equality (see [7, p.189])

$$(A_\lambda u_\lambda(t), u'_\lambda(t)) = \frac{d}{dt} \mathcal{J}_\lambda(t) \quad \text{a.e. } t \in]0, T[$$

where $\mathcal{J}_\lambda(t) = \frac{\lambda}{2} |A_H u(t)|_H^2 + \mathcal{J}((1 + \lambda A_H)^{-1} u(t))$, we denote by assumption $(\beta)_2$

$$\begin{aligned} \frac{\alpha_N}{2} \int_0^t |u'_\lambda(\tau)|_H^2 d\tau + \mathcal{J}_\lambda(u_\lambda(t)) &\leq \frac{1}{2\alpha_N} \|f\|_{L^2(Q_t)}^2 + \mathcal{J}_\lambda(u_0) + \\ &+ \sum_{i=1}^k \int_0^t \int_{\Omega_\lambda^i} g_{\lambda,i}^N(x, u_\lambda(\tau, x)) u'_\lambda(\tau, x) dx d\tau, \end{aligned} \quad (2.44)$$

where $u_0 = B^{-1}v_0$ (it is easy to see that $u_0 = u_\lambda(0)$). As $u_\lambda, u_0, \lambda^{\frac{1}{2}} g_{\lambda,i}^N$ and $\lambda^{-\frac{1}{2}} \text{mes } \Omega_\lambda^i$ are bounded, it follows that

$$\begin{aligned} \left| \int_0^t \int_{\Omega} g_{\lambda,i}^N(x, u_\lambda(\tau, x)) u'_\lambda(\tau, x) dx d\tau \right| &= \\ &= \lambda^{-\frac{1}{2}} \left| \int_{\Omega_\lambda^i} \int_0^t \frac{\partial}{\partial \tau} \int_0^{\lambda^{\frac{1}{2}} u_\lambda(\tau, x)} \lambda^{\frac{1}{2}} g_{\lambda,i}^N(\tau, \xi) d\xi d\tau dx \right| = \\ &= \lambda^{-\frac{1}{2}} \left| \int_{\Omega_\lambda^i} \int_{u_0(x)}^{u_\lambda(t, x)} \lambda^{\frac{1}{2}} g_{\lambda,i}^N(x, \xi) d\xi dx \right| \leq C. \end{aligned} \quad (2.45)$$

Since $v_0 \in W \cap L^\infty(\Omega)$ and β^{-1} is locally Lipschitz, we infer that $u_0 \in W \cap L^\infty(\Omega)$. By virtue of Lemma 2.2, it follows that $u_0 \in L^\infty(S)$. From this, we deduce that $j(u_0) \in L^1(S)$, and therefore, $u_0 \in D(\mathcal{J})$. Using the well-known inequality (see, e.g., [7, p.57])

$$\mathcal{J}_\lambda(u) \leq \mathcal{J}(u),$$

we deduce from (2.44) and (2.45) that

$$\frac{\alpha_N}{2} \int_0^t |u'_\lambda(\tau)|_H^2 d\tau + \frac{\lambda}{2} |A_H u_\lambda(t)|_H^2 + \frac{1}{2} |\text{grad } u_\lambda(t)|_H^2 \leq C. \quad (2.46)$$

Since $(1 + \lambda A_H)^{-1}$ is nonexpansive on H , it follows that

$$|y(t)|_H \leq C, \quad (2.47)$$

$$|y'(t)|_H \leq |u'(t)|_H. \quad (2.48)$$

From (2.46), (2.47) and (2.48) we get (2.40) and (2.41).

As $(1 + \lambda A_H)^{-1}$ is nonexpansive on $L^q(\Omega)$ for all $q \geq 2$, it follows that $\|y_\lambda(t)\|_{L^\infty(\Omega)} \leq N$. Noticing that $y_\lambda(t) \in W$ and using Lemma 2.2, we obtain (2.42).

Finally, we shall prove (2.43). Since $-\frac{\partial y_\lambda(t)}{\partial n} \in \sigma(y_\lambda(t))$, it follows from (2.42) and assumption (σ) that

$$\left\| \frac{\partial y_\lambda(t)}{\partial n} \right\|_{L^2(S)} \leq c. \quad (2.49)$$

Thus, from the equality

$$\|A_H y_\lambda(t)\|_{W'} = \sup_{\|\psi\|_W=1} (A_H y_\lambda(t), \psi) = \sup_{\|\psi\|_W=1} \left(\int_\Omega \text{grad } y_\lambda(t) \cdot \text{grad } \psi \, dx - \int_S \frac{\partial y_\lambda(t)}{\partial n} \psi \, dS \right)$$

we deduce immediately (2.43) as claimed.

3. PROOF OF THEOREM 1.1

By virtue of Lemma 2.3, it follows that

$$\{y_\lambda\} \text{ is bounded in } L^\infty(0, T; W), \quad (3.1)$$

$$\{y'_\lambda\} \text{ is bounded in } L^2(0, T; H), \quad (3.2)$$

$$\|y_\lambda(t) - u_\lambda(t)\|_H^2 \leq c\lambda, \quad (3.3)$$

$$\{A_\lambda u_\lambda\} = \{A_H y_\lambda\} \text{ is bounded in } L^\infty(0, T; W'). \quad (3.4)$$

By the Arzelà-Ascoli theorem, we obtain that on some subsequence convergent to zero of λ (for simplicity denoted again by λ) we have

$$y_\lambda \rightarrow u \text{ strongly in } C([0, T]; H). \quad (3.5)$$

It is well known that (3.1) (3.4) imply that

$$y_\lambda \rightarrow u \text{ weak-star in } L^\infty(0, T; W), \quad (3.6)$$

$$A_\lambda u_\lambda = A y_\lambda \rightarrow w \text{ weak-star in } L^\infty(0, T; W') \quad (3.7)$$

respectively. From (3.3) and (3.5) it follows that

$$u_\lambda \rightarrow u \text{ strongly in } C([0, T]; H). \quad (3.8)$$

Combining (3.8) and the boundness of v_λ , we infer that (on some subsequence of λ)

$$v_\lambda \rightarrow v \text{ weakly in } L^2(0, T; H), \quad (3.9)$$

$$v(t) \in Bu(t) \quad \text{a.e. } t \in]0, T[. \quad (3.10)$$

Now we shall prove that

$$y_\lambda \rightarrow u \text{ strongly in } C([0, T]; L^2(S)). \quad (3.11)$$

It is well known that the embedding operator from W into $L^2(S)$.

Therefore, from (3.1) it follows that for each $t \in [0, T]$,

$$\{y_\lambda(t)|_S\} \text{ is compact in } L^2(S). \quad (3.12)$$

Using the inequality (see [9, p.p. 47-49])

$$\|u\|_{L^2(S)} \leq \delta |\text{grad } u|_H + C_\delta u_H \quad (3.13)$$

where $\delta > 0$ is arbitrary small, we deduce

$$\begin{aligned} \|y_\lambda(t) - y_\lambda(s)\|_{L^2(S)} &\leq \delta |\text{grad } y_\lambda(t) - \text{grad } y_\lambda(s)|_H + \\ &+ C_\delta \|y_\lambda(t) - y_\lambda(s)\|_H. \end{aligned} \quad (3.14)$$

By virtue of (2.40), (2.41), from (3.14) it follows that

$$\|y_\lambda(t) - y_\lambda(s)\|_{L^2(S)} \leq C(\delta + \sqrt{|t-s|} C_\delta) \quad (3.15)$$

for $\lambda > 0, t, s \in [0, T]$. This means that $\{y_\lambda\}$ is equicontinuous in $C([0, T]; L^2(S))$. Combining this and (3.12) according to Arzelà-Ascoli's theorem, we obtain (3.11).

By virtue of (3.11) we can prove that

$$w(t) \in Au(t) \quad \text{a.e. } t \in]0, T[. \quad (3.16)$$

Indeed, it follows from (3.7) that for each $\psi \in L^1(0, T; W)$,

$$\int_0^T (A_\lambda u_\lambda(t), \psi(t)) dt \rightarrow \int_0^T (w(t), \psi(t)) dt, \quad (3.17)$$

and therefore,

$$\begin{aligned} \int_0^T (\text{grad } y_\lambda(t), \text{grad } \psi(t)) dt + \int_0^T \int_S -\frac{\partial y_\lambda(t)}{\partial n} \psi(t) dS dt \\ \rightarrow \int_0^T (w(t), \psi(t)) dt \end{aligned} \quad (3.18)$$

From (3.5), (3.6), we may infer that for $\psi \in L^1(0, T; W)$

$$\begin{aligned} \int_0^T (y_\lambda(t), \psi(t)) dt &\rightarrow \int_0^T (u(t), \psi(t)) dt, \\ \int_0^T (y_\lambda(t), \psi(t)) dt + \int_0^T (\text{grad } y_\lambda(t), \text{grad } \psi(t)) dt \\ &\rightarrow \int_0^T (u(t), \psi(t)) dt + \int_0^T (\text{grad } u(t), \text{grad } \psi(t)) dt, \end{aligned}$$

and therefore

$$\int_0^T (\text{grad } y_\lambda(t), \text{grad } \psi(t)) dt \rightarrow \int_0^T (\text{grad } u(t), \text{grad } \psi(t)) dt. \quad (3.19)$$

Because $-\partial y_\lambda(t)/\partial n \in \sigma(y_\lambda(t))$, and $\{y_\lambda\}$ is bounded on S , from (3.11) it follows that there exists a subsequence of λ (again denoted by λ) such that

$$-\frac{\partial y_\lambda}{\partial n} \rightarrow z \quad \text{weakly in } L^2(\Sigma_T) \quad (3.20)$$

$$\text{and } z(t) \in \sigma(u(t)) \quad \text{a.e. } t \in]0, T[. \quad (3.21)$$

Combining (3.18), (3.19), (3.20), we obtain

$$\int_0^T (\text{grad } u(t), \text{grad } \psi(t)) dt + \int_0^T \int_S z(t) \psi(t) dS dt = \int_0^T (w(t), \psi(t)) dt \quad (3.22)$$

for each $\psi \in L^1(0, T; W)$. Since ψ was arbitrary, we have

$$(\text{grad } u(t), \text{grad } \psi) + \int_S z(t) \psi dS = (w(t), \psi) \quad (3.23)$$

for each $\psi \in W$ and a.e. $t \in]0, T[$, i.e. (3-16) holds.

Now we shall prove

$$G_\lambda u_\lambda \rightarrow Gu \quad \text{weak-star in } L^2(0, T; W'), \quad (3.24)$$

i.e.

$$\lim_{\lambda \rightarrow 0} \int_0^T (G_\lambda u_\lambda(t) - Gu(t), \psi(t)) dt = 0 \quad (3.25)$$

for each $\psi \in L^2(0, T; W)$.

Forst, we remind that Gu is ^{well} defined because $u \in L^\infty(0, T; W) \cap L^\infty(Q_T)$, and so $u \in L^\infty(\Sigma_T)$.

It is easy to see that

$$(Gu(t), \psi(t)) = \sum_{i=1}^k \int_{\Omega_\lambda^i} \lambda^{-\frac{1}{2}} g(\bar{x}^i, u(t, \bar{x}^i)) \psi(t, \bar{x}^i) n_i(x) dx, \quad (3.26)$$

and therefore,

$$\begin{aligned} |(G_\lambda u_\lambda(t) - Gu(t), \psi(t))| &= \left| \sum_{i=1}^k \int_{\Omega_\lambda^i} \lambda^{-\frac{1}{2}} [g(\bar{x}^i, u_\lambda(t, x)) \psi(t, x) - \right. \\ &\quad \left. - \psi(t, \bar{x}^i) g(\bar{x}^i, u(t, \bar{x}^i))] n_i(x) dx \right| \leq \sum_{i=1}^k \int_{\Omega_\lambda^i} \lambda^{-\frac{1}{2}} [L_N(|u_\lambda(t, x) - y_\lambda(t, x)| + \\ &\quad + |y_\lambda(t, x) - y_\lambda(t, \bar{x}^i)| + |y_\lambda(t, \bar{x}^i) - u(t, \bar{x}^i)|) |\psi(t, x)| + \\ &\quad + |g(\bar{x}^i, u(t, \bar{x}^i))| |\psi(t, x) - \psi(t, \bar{x}^i)|] n_i(x) dx. \end{aligned} \quad (3.27)$$

Using the following inequality

$$\|u(x) - u(\bar{x}^i)\|_{L^2(\Omega_\lambda^i)} \leq c \lambda^{\frac{1}{2}} \|\text{grad } u\|_{L^2(\Omega_\lambda^i)}, \quad \forall u \in W, \quad (3.28)$$

it follows from (3.27) that

$$\begin{aligned} |(G_\lambda u_\lambda(t) - Gu(t), \psi(t))| &\leq c \sum_{i=1}^k \left[L_N (\lambda^{\frac{1}{2}} |A_H y_\lambda(t)|_H + \|\text{grad } y_\lambda(t)\|_{L^2(\Omega_\lambda^i)}) \right. \\ &\quad \cdot \|\psi(t)\|_{L^2(\Omega_\lambda^i)} + M_N \|\text{grad } \psi\|_{L^2(\Omega_\lambda^i)} (\text{mes } \Omega_\lambda^i)^{\frac{1}{2}} \Big] + \\ &\quad + c \sum_{i=1}^k \lambda^{-\frac{1}{4}} \|y_\lambda(t) - u(t)\|_{L^2(S_i)} \|\psi(t)\|_{L^2(\Omega_\lambda^i)}. \end{aligned} \quad (3.29)$$

For $u \in W$ and $x \in \Omega_\lambda^i$, it follows by the inequality

$$|u(y)|^2 \leq |u(\bar{y}, h_i(\bar{y}))|^2 + 2 \int_{h_i(\bar{y})}^{h_i(\bar{y}) + \sqrt{\lambda}} |u(y)| \left| \frac{\partial u(y)}{\partial y_n} \right| dy_m$$

that

$$\int_{h_i(\bar{y})}^{h_i(\bar{y}) + \sqrt{\lambda}} |u(y)|^2 dy_m \leq \sqrt{\lambda} (|u(\bar{y}, h_i(\bar{y}))|^2 + 2 \int_{h_i(\bar{y})}^{h_i(\bar{y}) + \sqrt{\lambda}} |u(y)| \left| \frac{\partial u(y)}{\partial y_m} \right| dy_m).$$

Multiplying this inequality by $n_i(T_i^{-1}y)$ and integrating over D_i , we get

$$\|u\|_{L^2(\Omega_\lambda^i)}^2 \leq c \sqrt{\lambda} (\|u\|_{L^2(S_i)}^2 + \|u\|_W^2).$$

Using the inequality $\frac{1}{L^2(S)} \|u\|_{L^2(S)} \leq c \|u\|_W$, it follows that

$$\|u\|_{L^2(\Omega_\lambda^i)} \leq c \lambda^{\frac{1}{4}} \|u\|_W, \quad \forall u \in W. \quad (3.30)$$

By virtue of (3.30) and the obvious inequality

$$\lim_{\lambda \rightarrow 0} \int_0^T \|\psi(t)\|_{H^1(\Omega_\lambda^i)} dt = 0, \quad (3.31)$$

we deduce (3.25) from (3.11) and (3.29).

Letting $\lambda \rightarrow 0$ in the equation

$$v_\lambda(t) + \int_0^t (A_\lambda u_\lambda(\tau) - G_\lambda u_\lambda(\tau)) d\tau = v_0 + \int_0^t f(\tau) d\tau \quad (3.32)$$

which is equivalent to (2.14)-(2.15), we obtain

$$v(t) + \int_0^t (w(\tau) - Gu(\tau)) d\tau = v_0 + \int_0^t f(\tau) d\tau \quad (3.33)$$

where the limit is in the sense of weak-star convergence in $L^2(0, T; W')$.

Let $\phi \in C^1([0, T])$, $\phi(T) = 0$, $\psi \in W$. Multiplying (3.33) by $(-\phi_t \psi)$ and integrating over Q_T , we get (1.12) and (1.13). This completes the proof of THEOREM 1.1.

REMARK. The problem obtained by taking $f(t, x, u)$ instead of $f(t, x)$ in (1.1)-(1.3), has a local (i.e. for T sufficiently small) solution in sense of Theorem 1.1, if $f(t, x, u)$ is measurable as a function of $(t, x) \in Q_T$ for each $u \in \mathbb{R}^1$ continuous as a function of $u \in \mathbb{R}^1$ for a.e. $(t, x) \in Q_T$, and if in addition, for each $N > 0$, there exists $C_N > 0$ such that

$$|f(t, x, u)| \leq C_N$$

for a.e. $(t, x) \in Q_T$, and $|u| \leq N$. We can prove this conclusion with the method used in [10].

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