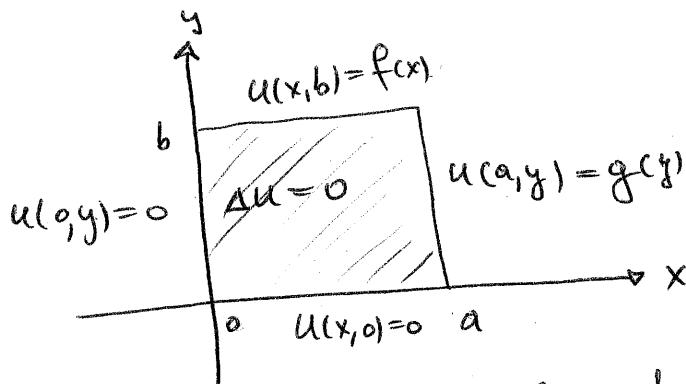


4.3. Potential in a rectangle

Example 1: $\Delta u = u_{xx} + u_{yy} = 0, 0 < x < a, 0 < y < b$
 $u(0, y) = 0, u(a, y) = g(y), 0 < y < b$
 $u(x, 0) = 0, u(x, b) = f(x), 0 < x < a$



Remark: We need to split this problem into two problems if we want to use separation of variables.

PDE1: $\Delta u_1 = 0$
 $u_1(0, y) = 0, u_1(a, y) = 0$ ← homogeneous vertical boundary
 $u_1(x, 0) = 0, u_1(x, b) = f(x)$

PDE2: $\Delta u_2 = 0$
 $u_2(0, y) = 0, u_2(a, y) = g(y)$
 $u_2(x, 0) = 0, u_2(x, b) = 0$ ← homogeneous horizontal boundary

We then have $u(x, y) = u_1(x, y) + u_2(x, y)$.
 We solve PDE1 using separation of variables: $u_1(x, y) = X(x)Y(y)$ and

get $X''Y + XY'' = 0 \Rightarrow \frac{X''}{X} = -\frac{Y''}{Y} = -\lambda^2$

Also $X(0) = 0, X(a) = 0, Y(0) = 0$.

From $X'' + \lambda^2 X = 0$ we find $X(x) = c_1 \sin(\lambda x) + c_2 \cos(\lambda x)$
 $X(0) = 0 \Rightarrow c_2 = 0, X(a) = 0 \Rightarrow \sin(\lambda a) = 0$ or $\lambda = \frac{n\pi}{a}, n=1, 2, \dots$
 $X_n(x) = \sin(\lambda_n x), \lambda_n = \frac{n\pi}{a}$.

$$Y'' - \lambda^2 Y = 0 \text{ gives } Y(y) = C_1 \sinh(\lambda y) + C_2 \cosh(\lambda y)$$

$$Y(0) = 0 \text{ so } C_2 = 0 \text{ and } Y(y) = \sinh(\lambda y)$$

$Y_n(y) = \sinh(\lambda_n y)$ and we have found

$$u_1(x, y) = \sum_{n=1}^{\infty} C_n X_n(x) Y_n(y) = \sum_{n=1}^{\infty} C_n \sinh\left(\frac{n\pi y}{a}\right) \sin\left(\frac{n\pi x}{a}\right)$$

From $u_1(x, b) = f(x)$ we get $\sum_{n=1}^{\infty} C_n \sinh\left(\frac{n\pi b}{a}\right) \sin\left(\frac{n\pi x}{a}\right) = f(x)$

Set $b_n = C_n \sinh\left(\frac{n\pi b}{a}\right)$ a constant.

Then $\sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{a}\right) = f(x)$ so $b_n = \frac{2}{a} \int_0^a f(x) \sin\left(\frac{n\pi x}{a}\right) dx$

and $C_n = \frac{b_n}{\sinh\left(\frac{n\pi b}{a}\right)}$. The solution of PDE1 is given by:

$$u_1(x, y) = \sum_{n=1}^{\infty} b_n \frac{\sinh\left(\frac{n\pi y}{a}\right)}{\sinh\left(\frac{n\pi b}{a}\right)} \sin\left(\frac{n\pi x}{a}\right), \quad b_n = \frac{2}{a} \int_0^a f(x) \sin\left(\frac{n\pi x}{a}\right) dx$$

Using a similar strategy we solve PDE2 to find $u_2(x, y)$:

$$u_2(x, y) = \sum_{n=1}^{\infty} C_n \sinh(\lambda_n x) \sin(\lambda_n y), \quad \lambda_n = \frac{n\pi}{b}, \quad n=1, 2, \dots$$

Remark: we swap $x \leftrightarrow y$ in PDE1 to get the solution for PDE2.

$$a \leftrightarrow b$$

$$f \leftrightarrow g$$

So one needs to pay attention on constants.

$$u_2(a, y) = g(y) = \sum_{n=1}^{\infty} C_n \sinh(\lambda_n a) \sin(\lambda_n y)$$

$$a_n = C_n \sinh(\lambda_n a) \text{ so } \sum_{n=1}^{\infty} a_n \sin\left(\frac{n\pi y}{b}\right) = g(y) \text{ gives } a_n = \frac{2}{b} \int_0^b g(y) \sin\left(\frac{n\pi y}{b}\right) dy$$

and $C_n = \frac{a_n}{\sinh\left(\frac{n\pi a}{b}\right)}$

by Fourier series

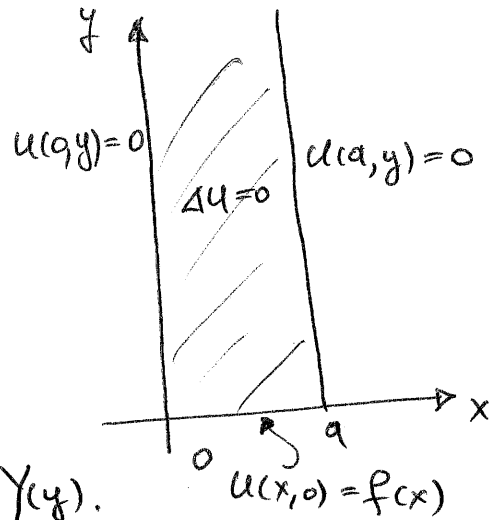
The general solution of PDE 2 is:

$$u_2(x, y) = \sum_{n=1}^{\infty} a_n \frac{\sinh\left(\frac{n\pi x}{b}\right)}{\sinh\left(\frac{n\pi a}{b}\right)} \sin\left(\frac{n\pi y}{b}\right), \text{ where}$$

$$a_n = \frac{2}{b} \int_0^b g(y) \sin\left(\frac{n\pi y}{b}\right) dy.$$

4.4. Potential in unbounded regions

Example 2: $\Delta u = 0, 0 < x < a, y > 0$
 $u(x, 0) = f(x), 0 < x < a$
 $u(0, y) = 0, u(a, y) = 0, 0 < y$
 $u(x, y)$ bounded as $y \rightarrow \infty$



We set up separation of variables $u(x, y) = X(x)Y(y)$.
 Then $\frac{X''}{X} = -\frac{Y''}{Y} = -\lambda^2$ (we need to set it to $-\lambda^2$, not λ^2 , because otherwise $Y(y)$ will become ~~unbounded~~ zero.)

$$X(0) = X(a) = 0$$

$$X'' + \lambda^2 X = 0 \text{ and } Y'' - \lambda^2 Y = 0$$

We find $X_n(x) = \sin(\lambda_n x), \lambda_n = \frac{n\pi}{a}, n=1, 2, \dots$

and $Y(y) = C_1 e^{-\lambda y} + C_2 e^{\lambda y}$ (It is more convenient to use this notation instead of $C_1 \sinh(\lambda y) + C_2 \cosh(\lambda y)$.)

$Y(y)$ bounded as $y \rightarrow \infty$ means that $C_2 = 0$ so $Y(y) = C_1 e^{-\lambda y}$. We found $Y_n(y) = e^{-\lambda_n y}, \lambda_n = \frac{n\pi}{a}$.

Putting all together we find:

$$u(x, y) = \sum_{n=1}^{\infty} C_n X_n(x) Y_n(y) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{a}\right) e^{-\frac{n\pi y}{a}}$$

From $u(x, 0) = f(x) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi x}{a}\right)$ we get $C_n = \frac{2}{a} \int_0^a f(x) \sin\left(\frac{n\pi x}{a}\right) dx$

Example 3:

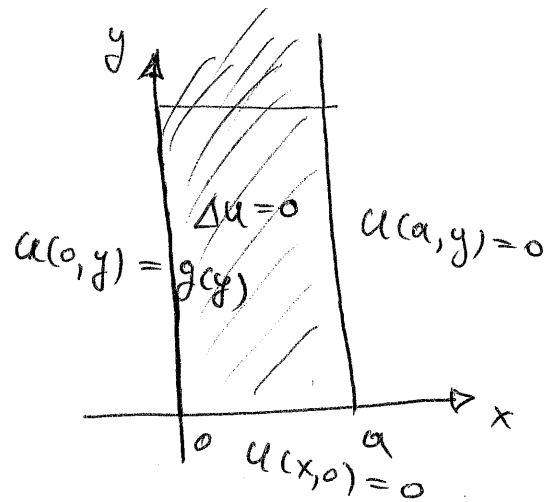
$$\Delta u = 0, 0 < x < a, 0 < y$$

$$u(x, 0) = 0, 0 < x < a$$

$$u(0, y) = g(y), 0 < y$$

$$u(a, y) = 0, 0 < y$$

$$u(x, y) \text{ bounded as } y \rightarrow \infty$$



$$u(x, y) = X(x) Y(y) \Rightarrow \frac{X''}{X} = -\frac{Y''}{Y} = \lambda^2$$

and:

$$X(a) = 0$$

$$Y(0) = 0$$

here we set it equal to λ^2 because if $-\frac{Y''}{Y} = -\lambda^2$ then $Y'' - \lambda^2 Y = 0$ and $Y(0) = 0$ will give $Y(y) = c_1 e^{\lambda y} + c_2 e^{-\lambda y}$, $c_2 = -c_1$
 $Y(y) = c_1 (e^{\lambda y} - e^{-\lambda y})$ but $Y(y)$ is bounded as $y \rightarrow \infty$ so $c_1 = 0$ and $Y \equiv 0$

$$\text{we find } \begin{cases} Y'' + \lambda^2 Y = 0 \\ X'' - \lambda^2 X = 0 \end{cases}$$

So $Y(y) = c_1 \cos(\lambda y) + c_2 \sin(\lambda y)$, $Y(0) = 0$ gives $c_1 = 0$

$Y(y) = \sin(\lambda y)$ (can take $c_1 = 1$ at this step)

$$X(x) = c_1 e^{-\lambda x} + c_2 e^{\lambda x} \quad \text{or}$$

$$X(x) = A \sinh(\lambda x) + B \cosh(\lambda x)$$

← this is a more convenient notation in this problem.

$X(a) = 0$ gives $A \sinh(\lambda a) + B \cosh(\lambda a) = 0$ and $A = -B \frac{\cosh(\lambda a)}{\sinh(\lambda a)}$

$$\text{so } X(x) = B \left(-\frac{\cosh(\lambda a) \sinh(\lambda x)}{\sinh(\lambda a)} + \cosh(\lambda x) \right)$$

$$= B \left(\frac{\cosh(\lambda x) \sinh(\lambda a) - \sinh(\lambda x) \cosh(\lambda a)}{\sinh(\lambda a)} \right)$$

$$= B \left(\frac{\sinh((a-x)\lambda)}{\sinh(\lambda a)} \right) \quad \left| \begin{array}{l} \text{we have used the identity:} \\ \sinh(a \pm b) = \sinh a \cosh b \pm \sinh b \cosh a \end{array} \right.$$

The constant B can depend on λ .

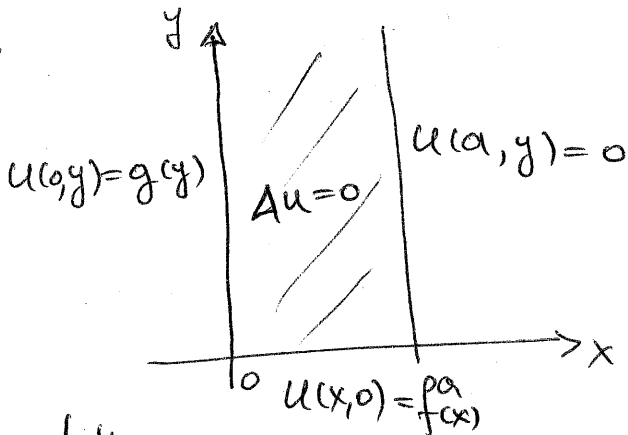
So

$$u(x, y) = \int_0^{\infty} B(\lambda) \frac{\sinh((a-x)\lambda)}{\sinh(\lambda a)} \sin(\lambda y) d\lambda$$

From $u(0, y) = g(y)$ we find $u(0, y) = \int_0^{\infty} B(\lambda) \sin(\lambda y) d\lambda$

so $B(\lambda) = \frac{2}{\pi} \int_0^{\infty} g(y) \sin(\lambda y) dy$. (Review ch. 2.10 and 1.9).

Example 4:



The solution to this PDE is Example 2 + Example 3.

4.5. Potential in a disk

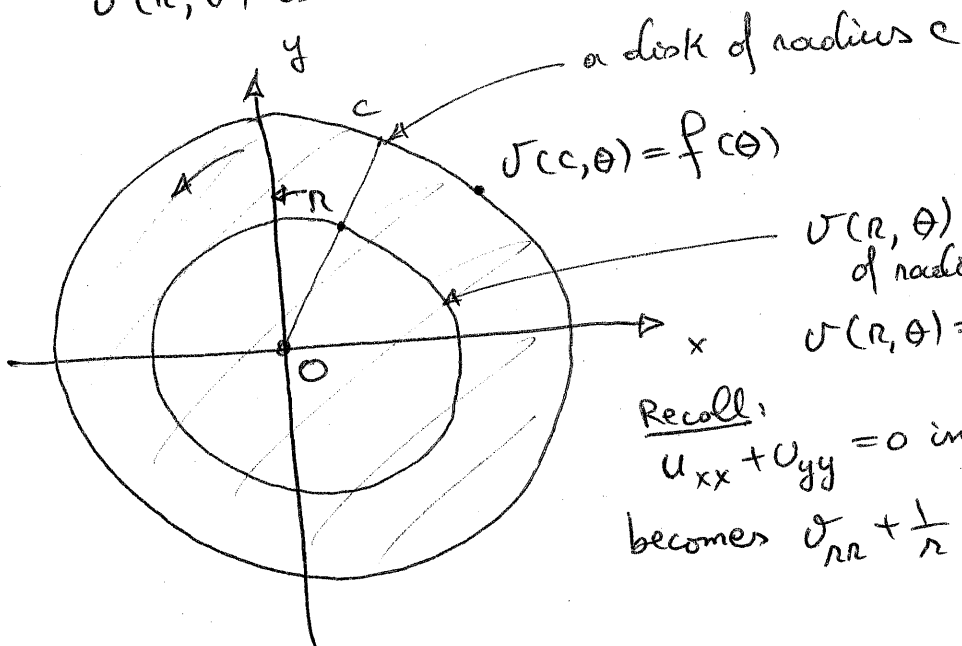
Example 5:

$$\nabla_{RR}^2 + \frac{1}{r} \nabla_r + \frac{1}{r^2} \nabla_{\theta\theta} = 0, \quad 0 \leq r < c$$

$$v(c, \theta) = f(\theta), \quad -\pi < \theta < \pi$$

$$v(r, \theta + 2\pi) = v(r, \theta), \quad 0 < r < c$$

$v(r, \theta)$ is bounded as $r \rightarrow 0$.



Recall:
 $u_{xx} + u_{yy} = 0$ in polar coordinates becomes $\nabla_{RR}^2 + \frac{1}{r} \nabla_r + \frac{1}{r^2} \nabla_{\theta\theta} = 0$.

Set $\psi(r, \theta) = R(r)\Theta(\theta)$ so

$$R''\Theta + \frac{1}{r}R'\Theta + \frac{1}{r^2}R\Theta'' = 0 \text{ gives}$$

$$(R'' + \frac{1}{r}R')\Theta = -\frac{1}{r^2}R\Theta'' \text{ or } \frac{R'' + \frac{1}{r}R'}{\frac{1}{r^2}R} = -\frac{\Theta''}{\Theta}$$

$$\text{or } \frac{r^2 R'' + r R'}{R} = -\frac{\Theta''}{\Theta} = \lambda^2 \left(\begin{array}{l} \text{if we set } = -\lambda^2 \text{ then } \Theta \text{ would} \\ \text{be exponential so not periodic;} \\ \text{we need } \Theta \text{ to be periodic} \end{array} \right)$$

$$\Theta'' + \lambda^2 \Theta = 0, \Theta(\theta + 2\pi) = \Theta(\theta)$$

so $\Theta(\theta) = A \cos(\lambda\theta) + B \sin(\lambda\theta)$ if this is periodic of period 2π then λ is an integer, so $\lambda = n, n = 0, 1, \dots$

if $\lambda = 0$ we get $\Theta_0(\theta) = \text{constant}$ so we pick $\Theta_0(\theta) = 1$.

if $\lambda = n, n = 1, 2, \dots$ $\Theta_n(\theta) = A_n \cos(n\theta) + B_n \sin(n\theta), n = 1, 2, \dots$

The equation for R is $r^2 R'' + r R' - n^2 R = 0, n = 0, 1, 2, \dots$

There are 2 cases:

① if $n = 0$ then $r^2 R'' + r R' = 0$ so $\frac{R''}{R'} = -\frac{1}{r}$ and so $R' = \frac{1}{r}$ or $R = \ln(r)$ if $R' \neq 0$

If $R' = 0$ then $R = \text{constant}$ so we take $R = 1$. Note that $R = \ln(r)$ does not work since $\lim_{r \rightarrow 0} R(r) = -\infty$.

$$R_0(r) = 1.$$

② $n \neq 0$. This is a Cauchy-Euler equation which cannot be solved by a characteristic equation. We know (and not prove) that the general solution is $R(r) = C_1 r^n + C_2 r^{-n}$. Now, since $R(r)$ is bounded as $r \rightarrow 0$ we must have $C_2 = 0$.

$$R_n(r) = r^n.$$

The fundamental solutions for this problem are

$$1, r^n \cos(n\theta), r^n \sin(n\theta)$$

or 1 and $A_n r^n \cos(n\theta) + B_n r^n \sin(n\theta)$ but they are the same.

$$\psi(r, \theta) = a_0 + \sum_{n=1}^{\infty} A_n r^n \cos(n\theta) + B_n r^n \sin(n\theta)$$

from the initial condition $v(c, \theta) = f(\theta)$ we find:

$$v(c, \theta) = a_0 + \sum_{n=1}^{\infty} A_n c^n \cos(n\theta) + B_n c^n \sin(n\theta) = f(\theta)$$

$$\text{so } a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(\theta) d\theta$$

$$A_n c^n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(\theta) \cos(n\theta) d\theta \quad \text{so } A_n = \frac{1}{\pi c^n} \int_{-\pi}^{\pi} f(\theta) \cos(n\theta) d\theta$$

$$B_n c^n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(\theta) \sin(n\theta) d\theta \quad \text{so } B_n = \frac{1}{\pi c^n} \int_{-\pi}^{\pi} f(\theta) \sin(n\theta) d\theta$$

using Fourier series.